Resonance assignment of ¹³C/¹⁵N labeled solid proteins by two- and three-dimensional magic-angle-spinning NMR

Mei Hong*

Department of Polymer Science & Engineering, University of Massachusetts, Amherst, MA 01003, U.S.A.

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Abstract

The comprehensive structure determination of isotopically labeled proteins by solid-state NMR requires sequence-specific assignment of 13 C and 15 N spectra. We describe several 2D and 3D MAS correlation techniques for resonance assignment and apply them, at 7.0 Tesla, to 13 C and 15 N labeled ubiquitin to examine the extent of resonance assignments in the solid state. Both interresidue and intraresidue assignments of the 13 C and 15 N resonances are addressed. The interresidue assignment was carried out by an N(CO)CA technique, which yields N_i-C α_{i-1} connectivities in protein backbones via two steps of dipolar-mediated coherence transfer. The intraresidue connectivities were obtained from a new 3D NCACB technique, which utilizes the well resolved C β chemical shift to distinguish the different amino acids. Additional amino acid type assignment was provided by a 13 C spin diffusion experiment, which exhibits 13 C spin pairs as off-diagonal intensities in the 2D spectrum. To better resolve carbons with similar chemical shifts, we also performed a dipolar-mediated INADEQUATE experiment. By cross-referencing these spectra and exploiting the selective and extensive 13 C labeling approach, we assigned 25% of the amino acids in ubiquitin sequence-specifically and 47% of the residues to the amino acid types. The sensitivity and resolution of these experiments are evaluated, especially in the context of the selective and extensive 13 C labeling approach.

Introduction

The determination of the three-dimensional structure of noncrystalline biological solids such as membrane proteins and insoluble protein aggregates by NMR remains an important and challenging problem today. To achieve this goal, it is desirable to have a repertoire of solid-state NMR techniques that are capable of yielding a large number of structural constraints from each experiment and sample. A necessary element in such a methodology is isotopic labeling, used to enhance the signal-to-noise ratios of the NMR spectra. Although site-specific ¹³C and ¹⁵N labeling has been shown to allow the measurement of structural constraints with high precision (Creuzet et al., 1991;

Shon et al., 1991; Ketchem et al., 1993; McDowell and Schaefer, 1996; Weliky and Tycko, 1996; Feng et al., 1997; Long et al., 1998), it is an inefficient process, requiring major synthetic efforts and yielding only a limited amount of information per experiment. Therefore, increasing attention has recently been directed to uniform ¹³C/¹⁵N labeling of proteins and to the design of experiments that yield multiple distances or torsion-angle constraints. However, before structural measurements on these uniformly labeled proteins can be evaluated, their NMR spectra need to be assigned to the amino acid sequences of the proteins.

Resonance assignment of proteins in solution is currently accomplished by a number of multidimensional correlation techniques that encode the chemical shift frequencies of various spins during different time periods and utilize scalar couplings to transfer the coherence between these spins (Wagner and Wüthrich,

^{*}To whom correspondence should be addressed at: Department of Chemistry, Iowa State University, Ames, IA 50011, U.S.A. E-mail: mhong@iastate.edu

1982; Ernst et al., 1987; Clore and Gronenborn, 1991; Grzesiek and Bax, 1992a; Bax, 1994; Cavanagh et al., 1996). In principle, analogous correlation schemes can also be employed in designing multidimensional magic-angle spinning (MAS) experiments to achieve resonance assignment of solid proteins (Tycko, 1996). In the solid state, one can use not only scalar couplings but also dipolar couplings for coherence transfer (Baldus and Meier, 1996; Lesage et al., 1997; Hong and Griffin, 1998). In fact, due to its sharp distance dependence $(1/r^3)$, the dipolar interaction can distinguish one-bond connectivities from multiple-bond connectivities quite effectively. Under MAS, heteronuclear and homonuclear dipolar couplings can be reintroduced by various radio frequency (rf) sequences (Baldus et al., 1994; Bennett et al., 1994; Nielsen et al., 1994; Gregory et al., 1995; Sun et al., 1995; Griffin et al., 1998; Hohwy et al., 1998; Rienstra et al., 1998).

In this paper, we demonstrate four 2D and 3D MAS correlation techniques that together yield the resonance assignment of the 76-residue protein, ubiquitin, to a significant extent. We first describe a 2D N(CO)CA correlation technique that manifests sequential $C\alpha_{i-1}$ - N_i connectivities along the protein backbone. The technique is formally complementary to, but simpler than, the ¹⁵N-¹³C correlation experiment introduced recently (Hong and Griffin, 1998). A similar N(CO)CA experiment was recently published by Straus et al. (1998); however, the detailed implementations of the two experiments differ, with substantially higher sensitivities achieved by the current experimental design. Since more complete assignment requires better site resolution than is feasible by 2D spectroscopy alone, we introduce a 3D NCACB technique that allows amino acid type assignment by utilizing the C β chemical shifts. The ^{15}N - $^{13}C\alpha$ cross sections of the 3D spectrum display distinct peaks that were unresolved in the 2D 15N-13C spectra. On the other hand, the ${}^{13}\text{C}\alpha$ - ${}^{13}\text{C}\beta$ cross sections of the 3D NCACB spectrum can be compared to a 2D ¹³C homonuclear correlation spectrum, obtained via ¹H-driven spin diffusion, which reveals ¹³C connectivities as off-diagonal intensities. We show that additional amino acid type assignment from this experiment can be extracted based on the chemical shifts of carbons further down the side chains. To better resolve and assign carbons with similar chemical shifts, which appear close to the diagonal of the spin diffusion spectrum, we also employ a double-quantum INADEQUATE experiment.

These heteronuclear and homonuclear correlation experiments are demonstrated on two ubiquitin samples. Both are uniformly labeled in ¹⁵N, but while one is uniformly labeled in ¹³C, the other is selectively and extensively labeled in ¹³C (Hong, 1999a; Hong and Jakes, 1999). The selective ¹³C labeling was achieved by using a specifically ¹³C-labeled glycerol, [2-¹³C] glycerol, as the sole carbon source in the defined media for protein expression. By breaking the ¹³C spin network with unlabeled carbon sites, selective and extensive ¹³C labeling reduces the effects of one-bond ¹³C-¹³C scalar couplings (∼55 Hz) and the multi-spin dipolar couplings, thereby enhancing the resolution of the NMR spectra. As analyzed previously, the spectral linewidth is the critical factor determining the extent and uniqueness of resonance assignment (Tycko, 1996). We recently showed that due to the selective and extensive ¹³C labeling, relatively well resolved spectra of proteins can be obtained, even when the samples were unoriented and uncrystallized, and the experiments were carried out under moderate spinning speeds (\sim 7 kHz) and at a moderate field strength (7.0 Tesla) (Hong, 1999a; Hong and Jakes, 1999). We also determined the labeling levels of each labeled carbon semi-quantitatively, using solution ¹³C-¹H correlation NMR. Most of these labeling levels agree well with the theoretical prediction based on the biosynthetic pathways (Tables 1 and 2 in (Hong, 1999a)). In this paper, we address some of the unique aspects of resonance assignment on such a selectively and extensively ¹³C-labeled protein. We show that knowledge of the labeling patterns and labeling levels affords additional information on resonance assignment that would be difficult to extract from the spectra of fully ¹³C-labeled proteins.

Pulse sequences for resonance assignment

N(CO)CA 2D correlation

We showed recently that sequential connectivities between an amide nitrogen of residuei and the $C\alpha$ site of residue i-1 can be obtained using a time-dependent, dipolar-mediated $^{15}N^{-13}C$ correlation experiment (Hong and Griffin, 1998). At long dipolar coherence transfer times, weak two-bond N_i - $C\alpha_{i-1}$ cross peaks between two adjacent amino acid residues can be observed in conjunction with strong one-bond N_i - $C\alpha_i$ cross peaks. But the dominant intraresidue cross peaks are of little utility for sequence-specific assignment, thus an experiment that selectively removes

these N_i - $C\alpha_i$ resonances will simplify the spectra and provide better site resolution. The N(CO)CA experiment is designed to select sequential N_i - $C\alpha_{i-1}$ cross peaks via two steps of coherence transfer, each involving strong one-bond dipolar couplings. The first step relays the magnetization of amide ^{15}N of residue i to the directly bonded carbonyl carbon of residue i-1, while the second transfer connects the CO_{i-1} with $C\alpha_{i-1}$.

The pulse sequence is displayed in Figure 1a and the corresponding phase cycles are listed in Table 1. After cross polarization from ¹H to ¹³C, a train of rotor-synchronized 180° pulses is applied on the ¹³C and ¹⁵N channels to reintroduce the ¹³C-¹⁵N dipolar coupling. This REDOR mixing period (Gullion and Schaefer, 1989), together with a pair of simultaneous 90° pulses, converts the ¹³C magnetization (C_x) to ¹⁵N antiphase magnetization (C_zN_v), which then evolves during the t1 period. Next, the chemical-shift-modulated ¹⁵N transverse magnetization is reconverted to ¹³C single-quantum coherence by an identical ¹³C-¹⁵N mixing period. Before the end of the second REDOR period, a short C-H dipolar dephasing period of about 30 us is introduced by turning off the ¹H decoupler. This dephasing period destroys the magnetization of the protonated Cα sites while retaining most of the carbonyl intensities. As a result, the intraresidue N_i -C α_i cross peaks are removed from the final spectrum. The selected carbonyl polarization is subsequently transferred to the neighboring $C\alpha$ via $^{13}C^{-13}C$ dipolar coupling. We chose the CMR7 sequence (Rienstra et al., 1998) to reintroduce the homonuclear dipolar interaction because of its relative insensitivity to resonance offset and rf field inhomogeneity, but other homonuclear recoupling sequences with comparable performance can also be employed (Hohwy et al., 1998, 1999). By choosing a short C-C mixing time, we can selectively recouple the one-bond ¹³CO-¹³Cα spin pairs. During the t_2 period, the C α signals are detected along with the untransferred carbonyl magnetization. Due to the double-quantum nature of the CMR7 Hamiltonian, the Cα resonances exhibit opposite sign of intensities from the CO signals.

The application of the CMR7 sequence is made possible by the moderate spinning speed used, since the required sevenfold increase in the ¹³C rf field strength over the spinning speed imposes strong demands on the ¹³C as well as the ¹H rf field strengths. If higher spinning speeds are used, other recoupling sequences with smaller ratios of the rf field to the spin-

ning speed (Sun et al., 1995; Hohwy et al., 1999) or with small rf duty cycles (Bennett et al., 1992) will be more suitable.

During the experiment, the proton carrier frequency is set on resonance for the aliphatic protons, except during the ¹⁵N evolution period, when it is shifted 1.3 kHz downfield to be on resonance for the amide proton. A moderate resolution enhancement in the ¹⁵N dimension is observed due to this on-resonance ¹H decoupling.

NCACB 3D correlation

In the 2D ¹⁵N- ¹³C correlation experiment described previously (Hong, 1999a), the spectral resolution is limited by overlapping signals in the region (¹⁵N, ¹³C) = (110-130 ppm, 52-62 ppm). With the exception of Pro, Gly, Ser, Thr, Val, and Ala, which exhibit relatively distinct ¹⁵N or ¹³Cα chemical shifts, most amino acids have their ${}^{15}N_i$ - ${}^{13}C\alpha_i$ resonances within this region. To reduce resonance overlap, we incorporated a third dimension into the correlation scheme. Sidechain carbons, whose chemical shifts have larger dispersions than the $C\alpha$ and the carbonyl carbons, are good candidates for the additional dimension. The simplest experiment is an NCACB technique, where, after the ¹⁵N and ¹³Cα chemical shifts are encoded in two indirect dimensions, the magnetization is transferred to Cβ by a homonuclear dipolar recoupling sequence and detected during the t₃ period. The 3D pulse sequence, shown in Figure 1b, resembles that of the N(CO)CA experiment to a large extent. The only differences are the removal of the C-H dephasing period and the reduction of the ¹H-¹³C contact time at the beginning of the sequence. The latter serves to minimize the CO signals while retaining most of the Ca intensities. In this way, we suppress the N_i - CO_{i-1} - $C\alpha_{i-1}$ pathway, which simplifies the spectrum, reduces the necessary ω_2 spectral width, and minimizes the size of the 3D data set.

The 13 C carrier frequency is set in the middle of the $C\alpha$ and $C\beta$ chemical shift region during most of the experiment, except for the 13 C α evolution (t_2) period, when it is shifted by 3.1 kHz to be downfield from the $C\alpha$ region. This enables off-resonance detection of the aliphatic carbons in the ω_2 dimension, which avoids potential zero-frequency artifacts that would interfere with the real signals.

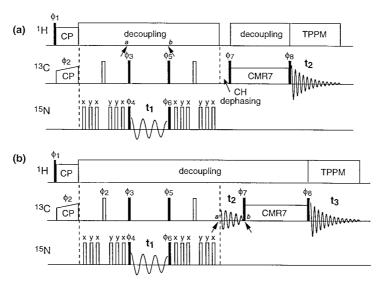


Figure 1. Pulse sequences for resonance assignments. (a) 2D N(CO)CA correlation. (b) 3D NCACB correlation. Open and filled rectangles indicate 180° and 90° pulses, respectively. ${}^{13}C^{-13}C$ double-quantum polarization transfer is achieved by the CMR7 sequence (Rienstra et al., 1998). Points a and b correspond to the frequency switching times on the ${}^{1}H$ (a) and ${}^{13}C$ (b) channel. The phase cycles are listed in Table 1.

Table 1. Phase cycles for the N(CO)CA and NCACB experiments. The phase numbering is indicated in Figure 1

φ1	ф2	ф3	ф4	ф5	Ф6	ф7	ф8	Receiver
+x	+x	+x	+ y	-x	+x-x	+y -y	+y -x -y +x	2301
$-\mathbf{x}$	+x	+x	+x	$-\mathbf{x}$	-x + x	+y-y	-x - y + x + y	3012
	+y	+y	+y	$-\mathbf{y}$	+y-y	-x + x	-y + x + y - x	0 1 2 3
	+y	+y	+y	$-\mathbf{y}$	-y + y	-x + x	+x +y -x -y	1230
	$-\mathbf{x}$	$-\mathbf{x}$	$-\mathbf{x}$	+x	-x + x	+y-y	+y-x-y+x	0123
	$-\mathbf{x}$	$-\mathbf{x}$	$-\mathbf{x}$	+x	+x-x	+y-y	-x - y + x + y	1230
	$-\mathbf{y}$	$-\mathbf{y}$	$-\mathbf{y}$	+y	-y + y	-x + x	-y + x + y - x	2 3 0 1
	-y	-y	-y	+ y	+y -y	-x + x	+x +y -x -y	3 0 1 2

Materials and methods

Sample preparation

The ubiquitin samples were obtained from VLI-research (Malvern, PA). The human ubiquitin gene was expressed in *E. coli*. Cells were grown in M9 minimal media, whose sole nitrogen source was (¹⁵NH₄)₂SO₄ and the sole carbon source was either [2-¹³C] glycerol (Cambridge Isotope Labs) or [U-¹³C₆] glucose. Below we shall refer to these two samples as ¹³C2-ubiquitin and U-¹³C, ¹⁵N-ubiquitin, respectively. For a typical ubiquitin expression experiment, a 5 L culture was induced and the solid paste was purified by column chromatography to homogeneity. The purity of the final products was analyzed by gel electrophoresis and amino acid composition

analysis to be >99%. The pools of purified protein solution were extensively dialysed against distilled water and lyophilized (http://www.vli-research.com/). The lyophilized protein was packed into a 4 mm MAS rotor and hydrated by direct addition of distilled and deionized H₂O. The water content was determined gravimetrically to be about 30% (w/w). It is well known that a hydrated powder yields sharper NMR lines than a lyophilized sample due to the reduction of structural heterogeneity (Poole and Finney, 1983; Kennedy and Bryant, 1990; Gregory et al., 1993). About 10 mg of U-¹³C,¹⁵N-ubiquitin was used for the N(CO)CA experiment, and about 30 mg of ¹³C2-ubiquitin was used for the 3D NCACB experiment and the ¹³C 2D experiments.

NMR spectroscopy

All NMR experiments were performed on a Bruker (Billerica, MA) DSX-300 spectrometer operating at 75.5 MHz for ¹³C and 30.4 MHz for ¹⁵N. A 4mm $^{1}H/^{13}C/^{15}N$ triple-resonance MAS probehead was used. The spinning speeds ranged from 6000 Hz to 7000 Hz, regulated by a Bruker MAS control unit. All experiments were conducted at room temperature, 293 \pm 2 K. Proton decoupling fields of 105–125 kHz were used. Typical carbon and nitrogen 90° pulse lengths were 4 µs and 6 µs, respectively. A linearly ramped rf field was applied on the ¹³C channel during ¹H-¹³C cross polarization to minimize dependence on the Hartmann-Hahn condition and maximize signal intensity (Metz et al., 1994). The CP contact time was typically 0.5 ms. During the acquisition period, TPPM decoupling (Bennett et al., 1995) was applied on the ¹H channel with a π pulse phase angle of $\pm 10^{\circ}$ to \pm 14°.

To set up the triple-resonance experiments, the efficiencies of homonuclear and heteronuclear polarization transfer were separately optimized before the two parts are combined in the final sequence. The efficiency of the REDOR-based (Gullion and Schaefer, 1989; Hing et al., 1992; Hong and Griffin, 1998) ¹⁵N-¹³C transfer was typically 20–25%, measured as the intensity of the ¹³C signal that passed the ¹⁵N-¹³C double- and zero-quantum filter relative to the intensity of the CP signal. The efficiency of the ¹³C-¹³C transfer under the CMR7 sequence was typically 40-45% after 1.1–1.2 ms of mixing. Higher efficiencies could not be attained, most likely due to the presence of 1.5–2 μ s windows between the 2π pulses, which were required to allow time for phase switching. For the N(CO)CA experiment, the C-H dipolar dephasing period, designed to destroy the signals of protonated carbons, reduced the intensity of the carbonyl resonance by another 20%. Thus, the overall efficiency of the N(CO)CA experiment, judged by the intensity of the negative Ca signals relative to their positive CP signals, varied from -5% to -10%. The NCACB experiment had a slightly higher overall efficiency, since no C-H dephasing period was necessary.

The signal-averaging time of the 3D NCACB experiment was minimized by carefully choosing the ω_1 and ω_2 spectral widths and the quadrature detection scheme. Both ^{15}N and $^{13}C\alpha$ chemical shifts were measured off-resonance so that only cosine data sets needed to be recorded. The dwell time for the ^{15}N (ω_1) dimension was chosen to be one rotor period, while the $^{13}C\alpha$ (t_2) evolution period was incremented at a step of

half a rotor period. At the spinning speed of 6000 kHz and the ^{15}N Larmor frequency of 30.4 MHz, the dwell time for ^{15}N evolution translates to an effective spectral width of 98.5 ppm. This is sufficient for covering most of the ^{15}N chemical shift range, except for the amine peaks, which were aliased into an empty region downfield from the amides. The recycle delay for the 3D experiment was optimized to 1.2 s. The relatively rapid ^{1}H T_{1} relaxation is a common feature of proteins due to the presence of various modes of motions.

The proton-driven ¹³C spin diffusion experiment was carried out using a 2D exchange sequence in which the ¹H decoupler was turned off during the mixing time (Szeverenyi et al., 1982). The cross peaks of one-bond and two-bond ¹³C-¹³C spin pairs were detected after a mixing time of 100 ms.

The 13 C dipolar INADEQUATE experiment was performed using CMR7 for homonuclear recoupling. The evolution period was incremented at a step of $^{1/7}$ of the CMR7 cycle time in order to generate a sufficiently large double-quantum spectral width (Hong, 1999b). At the spinning speed used (7 kHz), this corresponds to a ω_1 spectral width of 325 ppm.

Data processing

The spectra were processed using the XWINNMR software. A typical 2D spectrum was processed with 1024 and 256 complex points in the ω_2 and ω_1 dimensions, respectively. Gaussian broadening was employed in each dimension. For the 3D NCACB spectrum, 512, 128, and 128 points were used in the ω_3 , ω_2 and ω_1 dimensions, respectively.

Results and discussion

2D N(CO)CA spectrum of U-13C, 15N- ubiquitin

The N(CO)CA spectrum of uniformly 13 C, 15 N-labeled ubiquitin, shown in Figure 2a, demonstrates sequential assignment based on the N_i-C α_{i-1} cross peaks, in the absence of the intraresidue N_i-C α_{i} peaks. Effectively, it represents the difference spectrum between the one-bond 15 N- 13 C spectrum, acquired at short mixing times, and the multiple-bond 15 N- 13 C spectrum, acquired at long mixing times (Hong, 1999a). To illustrate this, we show the aliphatic region of the multiple-bond 15 N- 13 C spectrum in Figure 2b for comparison. The assignment proceeds by identifying resonances in the N(CO)CA spectrum whose (ω_1 , ω_2) frequencies uniquely equal the ω_1 frequency and

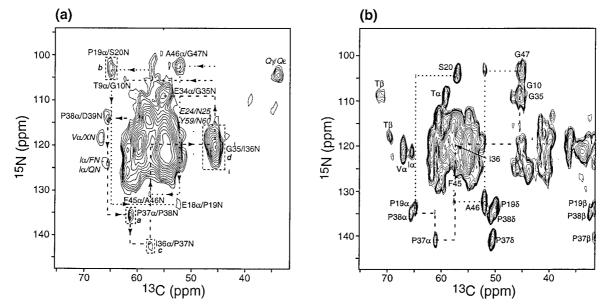


Figure 2. (a) 2D N(CO)CA spectrum of U- 13 C, 15 N-ubiquitin. (b) Multiple-bond 15 N- 13 C correlation spectrum of 13 C2-ubiquitin for comparison (Hong, 1999a). Four groups of sequential connectivities involving 14 resonances are assigned in the N(CO)CA spectrum. Linetypes differentiate the various connectivity paths while arrows guide the eye for these paths. A few resonances are assigned to the amino acid type or to one of a few possible sequence-specific residues. The acquisition time for the N(CO)CA spectrum was 13 h (number of scans per t_1 point = 96; number of t_1 points = 78; recycle delay = 3 s). Spinning speed: 6.5 kHz. 1 H decoupling field during acquisition: 109 kHz.

 ω_2 frequency of two separate peaks in the multiplebond ¹⁵N-¹³C spectrum. For example, the resonance labeled a at (61.3 ppm, 136.5 ppm) in the N(CO)CA spectrum (Figure 2a) corresponds to the 'junction' of resonances (61.1 ppm, 142.1 ppm) and (65.8 ppm, 136.3 ppm) in the multiple-bond ¹⁵N-¹³C spectrum (Figure 2b). These two 'parent' peaks have been previously assigned to prolines due to their uniquely downfield ¹⁵N chemical shifts. The presence of the cross peak in the N(CO)CA spectrum indicates that these two proline residues must be direct neighbors in the amino acid sequence. Thus the two parent proline peaks are unambiguously assigned to Pro37 and Pro38. Note that a weak cross peak between Pro37 and Pro38 is also observed in Figure 2b, due to direct through-space transfer from Pro38 ¹⁵N to Pro37 ¹³Cα. However, it is far less distinct than the corresponding peak in the N(CO)CA spectrum (Figure 2a). This demonstrates the efficiency of the two-step dipolar coherence transfer and the spectral simplification resulting from the removal of intraresidue resonances.

By superimposing the N(CO)CA spectrum with the multiple-bond 15 N- 13 C correlation spectrum, we can identify several other sequential connectivities. For example, resonance *b* at (65.0 ppm, 103.7 ppm) in Figure 2a links a Pro resonance (65.0 ppm, 134.5 ppm)

with a Ser resonance (57.0 ppm, 103.6 ppm) in Figure 2b. Since a unique Pro-Ser pair exists in the amino acid sequence of ubiquitin as Pro19-Ser20, this cross peak allows the assignment of the two intraresidue resonances in the multiplebond ¹⁵N-¹³C spectrum. In total, we assigned four groups of connectivities, indicated by different linetypes in Figure 2: Glu34–Gly35–Ile36–Pro37–Pro38– Asp39, Phe45-Ala46-Gly47, Glu18-Pro19-Ser20, and Thr9-Gly10. By design of the experiment, the consecutive turning points of each connectivity path alternate between finite and zero intensities. Those turning points with zero intensities correspond to the positions of the intraresidue ${}^{15}N_i$ - ${}^{13}C\alpha_i$ peaks that are manifested in the multiple-bond ¹⁵N-¹³C spectrum (Figure 2b) but have been removed by the N(CO)CA technique.

These resonances are assigned with slightly different degrees of confidence. For peaks that are not well resolved, the assignment particularly relies on the knowledge of the amino acid sequence of the protein. For example, the 13 C α chemical shift of Ile36 (57.4 ppm) is unmistakably defined by its cross peak with Pro37, resonance c (Figure 2a). Its 15 N chemical shift, on the other hand, is not clearly known from the N(CO)CA spectrum due to resonance overlap.

Since Ile36 neighbors Gly35, which has a characteristic 13 C α chemical shift well separated from all other amino acids, and since only one large resonance centered at (45.5 ppm, 120 ppm) (peak d) could result from glycine C α , we expect that the intraresidue 15 N- 13 C α resonance of Ile36 must appear at a 15 N frequency of about 120 ppm. This assignment is corroborated by the clear presence of the N-CO cross peak of Ile36 in the multiple-bond 15 N- 13 C spectrum (area not shown) (Hong, 1999a).

In addition to sequence-specific assignment, several more peaks in the N(CO)CA spectrum can be either assigned to amino acid types or to a limited number of sequence-specific possibilities, based on the unique chemical shift ranges of these amino acids. These inconclusive assignments, shown in italics in Figure 2a, are nevertheless useful in combination with other correlation experiments.

The resonance assignments were made independent of the liquid-state NMR chemical shifts known for ubiquitin, utilizing only the characteristic chemical shift ranges of amino acids and peptides and the connectivity patterns in the spectra. After the assignments had been made in the solid state, they were compared with the liquid-state chemical shifts of ubiquitin (Wang et al., 1995). For ambiguous assignments, the various possibilities are indicated in order not to exaggerate the extent of independent assignment by solid-state NMR.

Straus and co-workers recently demonstrated a similar N(CO)CA correlation experiment with a sequence that differs from the current version in three main aspects (Straus et al., 1998). First, the previous sequence employed CP instead of REDOR to produce polarization transfer from ¹⁵N to ¹³C. Second, it used a Gaussian 90° pulse instead of a transverse C-H dephasing period to select the carbonyl signals over Cα. Third, the previous sequence relied on protondriven spin diffusion to transfer the polarization from CO_{i-1} to $C\alpha_{i-1}$. We find that the coherence transfer schemes used in the current experiment yielded overall higher sensitivity than the previous sequence. This is judged from the signal-averaging times for the spectra, given similar sample amounts (\sim 10 mg) and similar experimental conditions. Cross peaks such as Pro37α–Pro38N, Ile36α–Pro37N, Glu18α–Pro19N, and Pro19α-Ser20N are visible in the current spectrum with comparable signal-to-noise ratios as the previous spectrum. However, the current 2D N(CO)CA spectrum was acquired in 13 h, which is significantly shorter than the 5.3 days reported in the previous paper. This suggests that one or several of our sequence elements have higher coherence transfer efficiency. Specifically, the efficiency of double-REDOR ¹⁵N-¹³C mixing was typically 20–25%. Although this efficiency is far below unity and lower than several cross polarization schemes (Hediger et al., 1994; Baldus et al., 1996), the transfer was achieved in a relatively short time (about 1.2 ms), thus minimizing signal loss due to transverse relaxation. The CO-Cα polarization transfer was also more rapid and better defined by homonuclear dipolar recoupling (about 1 ms) than by proton-driven spin diffusion (about 100 ms).

More complete sequential assignment of the ¹³C and ¹⁵N resonances in the protein backbone requires improved resolution compared to that of the 2D N(CO)CA spectrum. Although the distinct chemical shifts of several special amino acids such as proline, serine, and glycine result in partial site resolution, the spectrum in the (110–130 ppm, 52–62 ppm) region still contains many unresolved resonances. This can be partially attributed to the severe line broadening caused by uniform ¹³C labeling of the protein. Using the glycine signals as a benchmark, the linewidths of the CH₂ resonances of the uniformly labeled ubiquitin are about 2.2 ppm, while the linewidths in the ¹⁵N-¹³C spectra of selectively labeled ubiquitin are about 1.3 ppm.

3D NCACB spectrum of ¹³C2- ubiquitin

Enhanced site resolution can be achieved either by reducing the spectral linewidths or by increasing the dimensionality of the NMR spectrum. Narrower lines may be obtained by crystallizing the protein, which is possible for the well-behaved ubiquitin (Cook et al., 1979). However, we consider such a strategy not particularly meaningful, since it cannot be applied to noncrystalline and disordered proteins, which are the ultimate goal of the current technique development. Higher-resolution spectra may also be achieved by increasing the spinning speed (e.g., to 30 kHz) and using very high magnetic field strengths (e.g., 17.6 Tesla). However, since these conditions are not yet routine, we have pursued the second route, three-dimensional spectroscopy, to enhance the site resolution of solid protein spectra. Specifically, Cβ chemical shifts are utilized to disperse the ^{15}N - $^{13}C\alpha$ resonances in a third dimension (Grzesiek and Bax, 1992b).

The [2- 13 C] glycerol labeling scheme generates 13 C α - 13 C β spin pairs for 9 out of the 20 amino acids, thus it is reasonably suited for demonstrating the 3D NCACB experiment. The 3D spectrum

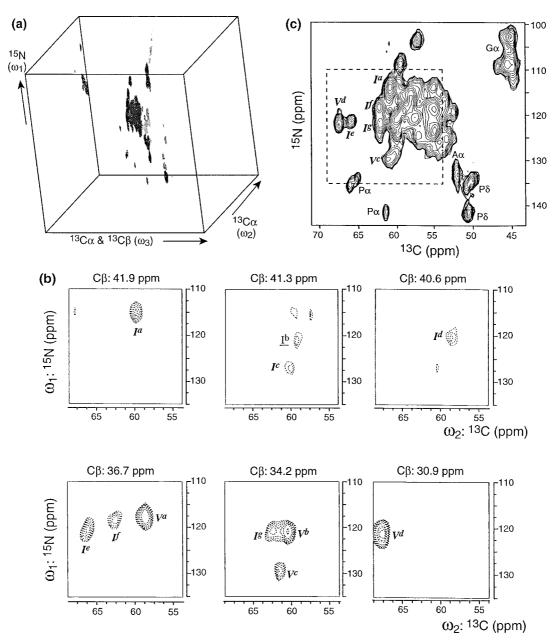


Figure 3. 3D NCACB spectrum of 13 C2-ubiquitin for amino acid type assignment. (a) 3D spectrum. Dark and light contours represent positive and negative intensities, respectively. (b) Selected 15 N_i- 13 C α_i (ω_1 - ω_2) cross sections from various 13 C β frequencies. Dashed contours indicate negative intensities. Bold italics indicate assignment confirmed by other experiments. Underlines indicate tentative assignment. (c) 2D one-bond 15 N- 13 C correlation spectrum for comparison (Hong, 1999a). The area enclosed by a dashed square corresponds to the spectral region shown in the cross sections. Note the higher resolution of the NCACB spectrum due to dispersion along the third dimension. The acquisition time for the 3D experiment was 52 h (number of scans per t_1 point = 32; number of t_1 points = 40; number of t_2 points = 116; recycle delay = 1.2 s). Spinning speed: 6 kHz. 11 H decoupling field during acquisition: 111 kHz.

of ¹³C2-ubiquitin is shown in Figure 3a. The positive intensities (dark contours) correspond to the Ca resonances while the negative intensities (light contours) correspond to the Cβ peaks. Figure 3b displays six $^{15}N^{-13}C\alpha$ (ω_1 - ω_2) cross sections extracted from various Cβ (ω₃) frequencies. The assignment begins with the identification of all four valine residues (labeled as V^a–V^d) in ubiquitin. These manifest as intense peaks with large round lineshapes in the region (13 C α , 13 C β) = (58–68 ppm, 30–37 ppm). The assignment is made on two grounds. First, valine is the only amino acid with 100% labeling for both Ca and C β in the [2-¹³C] glycerol labeling scheme. Thus the $^{13}\text{C}\alpha$ - $^{13}\text{C}\beta$ spin-pair probability of valine is the highest among all 20 amino acids, and the corresponding signal intensities should be the largest. Second, the average valine Cα and Cβ chemical shifts are 61 ppm and 31 ppm, respectively (Wüthrich, 1986). These are consistent with the frequencies of the four resonances when conformation-dependent secondary chemical shifts (Wishart et al., 1992) are taken into account. One of these four valine resonances, V^d, is corroborated by the time-dependent ¹⁵N-¹³C spectra as the resonance at $(^{15}N, ^{13}C) = (122.0 \text{ ppm},$ 67.2 ppm) (Hong, 1999a). The other three valine peaks had not been assigned previously in the multiplebond 2D ¹⁵N-¹³C spectrum (Figure 2b) due to insufficient spectral resolution. This demonstrates that 3D NCACB spectroscopy significantly surpasses 2D ¹⁵N-¹³C spectroscopy in the information content.

The amino acid with the next highest ${}^{13}\text{C}\alpha$ - ${}^{13}\text{C}\beta$ spin pair probability is isoleucine: its Cβ is 100% labeled while its $C\alpha$ is labeled at 50%, thus the $C\alpha$ - $C\beta$ pairs are doubly labeled at \sim 50%. In the remaining resonances in the ¹⁵N-¹³Cα cross sections, the peaks labeled I^a (115.1 ppm, 59.8 ppm, 41.9 ppm), I^c (127.4 ppm, 60.0 ppm, 41.3 ppm), I^d (120.4 ppm, 58.5 ppm, 40.6 ppm), Ie (120.4 ppm, 66.3 ppm, 36.7 ppm), and I^t (118.5 ppm, 62.5 ppm, 36.7 ppm) can be assigned to isoleucine with confidence, since they show relatively strong intensities and fall into the characteristic chemical shift regions of isoleucine, which are centered around (123 ppm, 60 ppm, 37 ppm) (Wüthrich, 1986). In addition, I^a, I^e and I^f have clear correspondences in the one-bond ¹⁵N-¹³C 2D spectrum (Figure 3c) (Hong, 1999a), while I^c and I^d are confirmed by the ¹³C spin diffusion spectrum (Figure 4) and the INADEQUATE spectrum (Figure 5) described below.

The combined use of signal intensities based on the known ¹³C labeling levels and the characteristic chem-

ical shift ranges reduces the ambiguity of the amino acid type assignment. For example, the Cβ chemical shifts of valine range from 29 ppm (α-helix) to 38 ppm $(\beta$ -sheet), while the C β chemical shifts of isoleucine range from 35 ppm (α -helix) to 44 ppm (β -sheet) (Spera and Bax, 1991). The partial overlap would have made the assignment of the V^a (¹³Cβ: 36.7 ppm) peak in Figure 3b uncertain. However, the distinctly stronger intensity of this resonance indicates that it must be valine. Additional chemical shift correlation also aids the assignment. For example, the resonance labeled Ig has an intensity and chemical shift values that could make it a possible valine if the NCACB spectrum alone is considered. However, in the ¹³C spin diffusion spectrum (Figure 4), the corresponding Cα-Cβ peak at (60.8 ppm, 34.5 ppm) shares a cross peak with an upfield methyl signal at 9.5 ppm. Since valine Cγ is unlabeled and resonates around 19 ppm, while isoleucine Cδ1 is labeled and resonates around 11 ppm, this Cα-Cβ resonance belongs unequivocally to an isoleucine. When all spectra are evaluated, only one resonance in Figure 3b, I^b, cannot be conclusively assigned.

The 3D NCACB spectrum was acquired in a relatively short time: 32 scans were coadded for each slice, resulting in slightly over two days for the entire experiment. With this signal-averaging time, amino acids derived from the citric acid cycle, which have lower levels of doubly labeled $C\alpha$ - $C\beta$ pairs than valine and isoleucine, could not be observed clearly. According to the labeling pattern for [2-¹³C] glycerol (Table 2 in (Hong, 1999a)), the Cα and Cβ sites are labeled at 25% and 50%, respectively, for Glu, Gln, Arg, and Pro, which derive from α -ketoglutarate. For Asp, Asn, Met, and Thr, whose common precursor is oxaloacetate, the Ca and CB sites are labeled at 50% and 25%, respectively. Therefore, the probability for doubly labeled $C\alpha$ - $C\beta$ pairs in these amino acids is about 12.5%, which is far below the 100% for valine and the 50% for isoleucine.

Although uniform 13 C labeling in principle allows more resonances to be observed in the 3D spectrum, it will also deteriorate the spectral resolution due to the 13 C- 13 C J-coupling and dipolar couplings. The reduced resolution will not only complicate the analysis of the 15 N- 13 C α cross sections but will also make it difficult to separate the 13 C β resonances, thus hindering the extraction of the 2D 15 N- 13 C α cross sections in the first place.

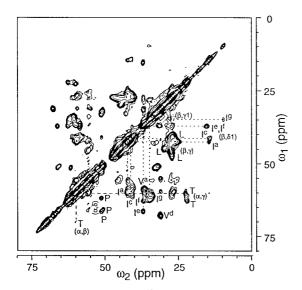


Figure 4. 2D proton-driven ^{13}C spin diffusion spectrum of $^{13}\text{C2}$ -ubiquitin, acquired with a mixing time of 100 ms. Greek letters in brackets specify the type of cross peaks. Superscripts for valine and isoleucine match the assignment in the NCACB spectrum. Dotted lines guide the eye for the connectivity paths. The acquisition time was 15 h (number of scans per t_1 point = 48; t_1 dwell time = 57.1 μ s; number of t_1 points = 400; recycle delay = 2.8 s). Spinning speed: 7 kHz. 1 H decoupling field during acquisition: 116 kHz.

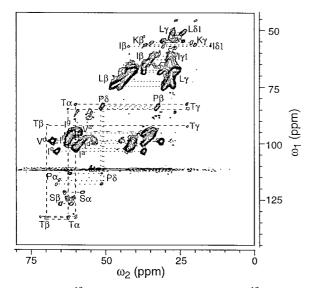


Figure 5. 2D 13 C dipolar INADEQUATE spectrum of 13 C2- ubiquitin, acquired with a double-quantum excitation time of 571 μs. Superscripts for valine and isoleucine match the assignment in the NCACB spectrum. Dotted lines guide the eye for the connectivity paths. The acquisition time was 22.5 h (number of scans per t_1 point = 64; t_1 dwell time = 40.8 μs; number of t_1 points = 800; recycle delay = 1.5 s). Spinning speed: 7 kHz. 1 H decoupling field during acquisition: 125 kHz.

 $2D^{13}C$ spin diffusion spectrum of $^{13}C2$ -ubiquitin

Amino acid type assignment can also be made by direct correlation of the ¹³C chemical shifts in a 2D ¹³C homonuclear correlation spectrum. Such a spectrum may be obtained in a number of ways, depending on the method of polarization transfer and the coherence order involved. We first show the result of a ¹H-driven spin diffusion experiment, which involves only single-quantum ¹³C coherence (Szeverenyi et al., 1982; McDermott et al., 1999). The mixing period for polarization transfer can be adjusted to probe connectivities between carbons separated by varying numbers of bonds.

In the spectrum of ¹³C2-ubiquitin acquired with 100 ms of mixing time (Figure 4), five types of residues, Ile, Val, Leu, Thr, and Pro, are identified with relative ease. Isoleucines are labeled at the C α , C β , C γ 1 and C δ 1 positions by the [2- 13 C] glycerol labeling scheme. Among these sites, Cy1 has the lowest labeling level and CB the highest, thus cross peaks between Cy1 and Cδ1 are not observed while cross peaks between Cβ and Cδ1 can be easily identified. The distinct upfield shift of Cδ1 (~11 ppm) allows the assignment of four peaks in the region (34–43 ppm, 10–16 ppm) to the Cβ-C δ 1 couplings of isoleucine. Once the Cβ chemical shifts are identified, Cβ-Cα cross peaks are found by tracing the connectivity paths. Five Cα-Cβ cross peaks at (60.8 ppm, 34.5 ppm), (62.8 ppm, 37.1 ppm), (66.3 ppm, 37.1 ppm), (59.5 ppm, 42.2 ppm), and (59.5 ppm, 42.2 ppm), agree with the I^g, I^f, I^e, I^c, and Ia resonances in the NCACB spectrum, respectively. The C β -C δ 1 peak of I^f at (36.7 ppm, 15.2 ppm) overlaps with that of Ie, which explains the particularly strong intensity of this resonance. Finally, three isoleucine Cβ-Cγ1 cross peaks are identified near the spectral diagonal in the region (34–37 ppm, 29-33 ppm).

Valine $C\alpha$ - $C\beta$ peaks can be identified based on their strong intensities due to 100% labeling and the distinctly downfield shift of the $C\alpha$ site. Thus two resonances at (67.8 ppm, 31.0 ppm) and (58.5 ppm, 36.7 ppm) are assigned to valine V^d and V^a , respectively. Leucines have 100% labeled $C\beta$ - $C\gamma$ pairs, whose unique chemical shifts around (40 ppm, 25 ppm) allow them to be assigned straightforwardly. There are nine leucine residues in ubiquitin, thus resonance overlap in this region is relatively severe, and only three major peaks can be resolved. These $C\beta$ - $C\gamma$ peaks are missing in the NCACB spectrum due to the long distance between the amide ^{15}N and the $C\beta$.

Threonines are characterized by strong $C\alpha$ - $C\gamma$ cross peaks (doubly labeled at ~25%) around (60 ppm, 20 ppm) and weak Cβ-Cα cross peaks (doubly labeled at \sim 12.5%) near the spectral diagonal around (70 ppm, 60 ppm). The intensity distribution of these peaks results from the low labeling level of C β ($\sim 25\%$), which is also reflected by the lack of C β -C γ cross peaks in the spectrum. Based on the strong $C\alpha$ - $C\gamma$ resonances, two threonine resonances are identified, out of a total of seven threonine residues in ubiquitin. Finally, the prolines are labeled at the $C\alpha$ (\sim 25%), CB (50%) and C δ (100%) sites of the ring. The directly bonded and \sim 25% labeled C α -C δ pairs give rise to three strong resonances in the region (60–66 ppm, 50– 52 ppm). These assignments corroborate the result of the ¹⁵N-¹³C 2D spectra.

The proton-driven spin diffusion experiment contains two 90° pulses before and after the mixing period. Thus it essentially constitutes a stimulatedecho experiment, which refocuses the magnetization at $t_2 = t_1$. According to Fourier transform theory, this indicates that the widths of the diagonal peaks correspond to the intrinsic homogeneous linewidths of the spectrum (Schmidt-Rohr and Spiess, 1994). The presence of sharp diagonal peaks with truncation wiggles indicates that the homogeneous linewidth of the spectrum is smaller than the typical linewidths of the off-diagonal peaks. This suggests possibilities for further improvement of the spectral resolution. Based on the non-diagonal intensities, the linewidths of methyl (e.g. Ile C δ 1), methylene (e.g., Leu C β), and methine (e.g., Val Cα) resonances are approximately 0.8 ppm, 1.5 ppm, and 1.2 ppm, respectively.

The initial rate of spin diffusion depends on the square of the dipolar coupling between the two exchanging spins and on an integral that describes the overlap between two ¹³C resonances in the presence of proton coupling (Linder et al., 1985). As a result, it is necessary to use hundreds of milliseconds or longer for polarization transfer over several angstroms. In this sense spin diffusion is an inefficient method for coherence transfer. Due to the complexity of the overlap integral, cross peak intensities in a spin diffusion spectrum do not correlate with internuclear distances quantitatively. Moreover, the spin diffusion spectrum suffers from a dominant diagonal, which obscures cross peaks close to it. To probe the ¹³C-¹³C bond connectivities more efficiently, with more emphasis on directly bonded carbons, and without neglecting spin pairs with small chemical shift differences, we employed a ¹³C double-quantum technique.

2D ¹³C dipolar-INADEQUATE spectrum of ¹³C2-ubiquitin

The dipolar-mediated INADEQUATE experiment correlates the double-quantum signals of two coupled spins with the single-quantum chemical shift. The double-quantum coherence can be excited by the ¹³C-¹³C dipolar coupling, which is reintroduced by the CMR7 sequence in our experiment (Hong, 1999b). A double-quantum excitation time of 571 µs, corresponding to one cycle of the sequence under 7 kHz of sample spinning, was employed. Figure 5 displays the aliphatic region of the INADEQUATE spectrum of ¹³C2-ubiquitin. Two groups of intense signals are observed. First, in the region $\omega_1 = 93-105$ ppm, the Cα-Cβ couplings of two valines (V^a and V^d) and five isoleucine residues (Ia, Id, Ie, If, and Ig) are observed. These corroborate the assignment of the spin diffusion spectrum and the 3D NCACB spectrum. In particular, the pair of signals at (98.7 ppm, 58.0 ppm) and (98.7 ppm, 40.7 ppm) corresponds to the I^d C α -C β resonance in the NCACB spectrum but was unresolved in the spin diffusion spectrum. This attests to the higher resolution of the INADEQUATE spectrum compared to the spin diffusion experiment. Second, in the region $\omega_1 = 60-80$ ppm, the strong and broad peaks can be assigned to the 100% labeled Cβ-Cγ pairs of leucine residues. The weaker and more closely spaced resonances are assigned to the Cβ-Cγ1 couplings of isoleucine based on the chemical shifts and the \sim 25% labeling level of the ¹³Cβ-¹³Cγ1 pairs (Hong, 1999a). These isoleucine resonances were easy to miss in the spin diffusion spectrum due to their proximity to the strong diagonal. The double-quantum filtration of uncoupled spins by the INADEQUATE technique simplifies the detection of these coupled spins with similar chemical shifts.

In addition to these intense resonances, three clusters of weak cross peaks are observed in the spectrum. From $\omega_1=113$ ppm to 135 ppm, we identified proline $C\alpha\text{-}C\delta$, serine $C\beta\text{-}C\alpha$, and threonine $C\beta\text{-}C\alpha$ resonances. The observation of the serine peaks is at first surprising, since only $C\alpha$ should be labeled in this amino acid according to the $[2^{-13}C]$ glycerol labeling protocol. However, since the $C\alpha$ site is labeled at 100%, even a low scrambling level of 15% at $C\beta$ would have increased the $^{13}C\beta\text{-}^{13}C\alpha$ spin-pair probability to 15%. This would come close to the spin-pair probability of many amino acids produced from the citric acid cycle, whose cross peaks are detectable in the spectrum. For example, the threonine $C\beta\text{-}C\alpha$ sites, which resonate slightly downfield from the ser-

ine peaks in the ω_1 dimension, are doubly labeled at about 12.5%. Therefore, the manifestation of the serine cross peaks indicates label scrambling at C β , which had not been detected in previous heteronucle ar and homonuclear spectra. It demonstrates the capability of double-quantum filtration in revealing low percentages of spin pairs.

From $\omega_1=80$ ppm to 93 ppm, proline C δ -C β , threonine C α -C γ and C β -C γ cross peaks are identified. The presence of a large splitting between (92.5 ppm, 70.0 ppm) and (92.5 ppm, 22.5 ppm), which is assigned to a threonine C β -C γ pair, attests to the broadband nature of the technique. Again, a strong correlation is found between the signal intensities and the labeling levels. For example, the threonine C β -C γ sites (\sim 12.5% labeling) exhibit much weaker signals than the C α -C γ pairs (\sim 25% labeling), although they are better recoupled due to the shorter internuclear distance.

In the most upfield region of the ω_1 dimension (45–57 ppm), the cluster of resonances is tentatively assigned to isoleucine C β -C δ 1, lysine C β -C γ , and leucine C γ -C δ 1 based on the chemical shifts. The assignment of the lysine and leucine peaks assumes scrambling at the C β site and the C δ 1 site, respectively.

Assignment from complementary $[1^{-13}C]$ glucose and $[2^{-13}C]$ glycerol labeling

Due to the partial complementarity of the [1-13C] glucose and [2-¹³C] glycerol labeling schemes, additional resonance assignment can be extracted by comparing the spectra of the differently labeled samples. For example, in the ¹⁵N-¹³C spectrum of ¹³C2-ubiquitin, two adjacent resonances at (121.7 ppm, 65.7 ppm) and (122.0 ppm, 67.2 ppm) were originally assigned to valine Ca and isoleucine Ca without distinguishing between the two (Hong, 1999a). In the ¹⁵N-¹³C spectrum of ¹³C1-ubiquitin (not shown), however, only one peak is observed at (122.0 ppm, 65.5 ppm). Since valine $C\alpha$ is unlabeled in the [1- 13 C] glucose scheme, while isoleucine $C\alpha$ is labeled, the peak in the ¹³C1-ubiquitin spectrum is definitively assigned to isoleucine N-Ca. Thus the remaining peak in the ¹³C2-ubiquitin spectrum is attributed to valine.

Summary of ^{13}C and ^{15}N assignment of ubiquitin by solid-state NMR

We have developed and demonstrated several 2D and 3D MAS correlation techniques to test the extent of ¹³C and ¹⁵N assignment of proteins in the solid state. These include three heteronuclear experiments:

MQIFV KTLTG
KTITL EVEPS
DTIEN VKAKI
QDKEG IPPDQ
QRLIF AGKQL
EDGRT LSDYN
IQKES TLHLV
LRLRG G

Figure 6. Amino acid sequence of ubiquitin. Underlined residues have been assigned by the techniques described here and in two recent publications (Hong, 1999a, b).

N(CO)CA, NCACB, and the time-dependent ¹⁵N-¹³C correlation described previously (Hong, 1999a), and two homonuclear techniques: proton-driven ¹³C spin diffusion and ¹³C dipolar INADEQUATE (Hong, 1999b). Combining all the spectra, we are able to assign the backbone ¹⁵N and ¹³C resonances and some of the side-chain carbons of 19 out of 76 residues in ubiquitin in a sequence-specific fashion. These conclusively assigned residues are underlined in the amino acid sequence of ubiquitin in Figure 6. In addition, a larger number of resonances has been assigned to the amino acid types, as compiled in Table 2. These assignments were made in the solid state based on the spectral connectivities, the characteristic chemical shifts, and the selective and extensive ¹³C labeling patterns. Since they were made independent of the solution-state chemical shifts, they represent the true extent of assignment currently feasible in our hands for a protein of this size.

The amino acid type assignment involved various side-chain 13 C chemical shifts, including C β , C γ , and Cδ. It can be seen that the most completely and easily assigned residues are either those with distinct chemical shifts, or those that are labeled at high levels and with high selectivity. Thus, all proline and valine resonances were type-assigned. Interestingly, although all three prolines in ubiquitin could be identified sequence-specifically, no valine resonances could be assigned sequentially, due to the lack of unique chemical shifts or unique labeling pattern of their neighboring residues. In addition to these two amino acids, several other amino acids lend themselves conveniently to type-assignment by virtue of their distinct chemical shifts: glycine (Cα), alanine (Cβ), leucine (C β and C γ), isoleucine (C δ 1), serine (C β), threonine $(C\beta)$, lysine $(C\epsilon)$, and arginine $(C\delta)$ (Tycko, 1996).

Table 2. Summary of type-assigned amino acid residues in ubiquitin, using ¹⁵N-¹³C and ¹³C-¹³C MAS multidimensional correlation spectroscopy

Amino acid	Total no.	No. assigned	Amino acid	Total no.	No. assigned
Gly	6	3a,c	Gln	6	0
Ser	3	3a-c,f	Pro	3	$3^{a-c,e-g}$
Ala	2	2^{a-c}	Arg	4	$1^{a,b,g}$
Leu	9	$3^{b,e-g}$	Asp	5	0
Val	4	$4^{a-e,f}$	Asn	2	$1^{b,g}$
His	1	1 ^h	Met	1	$1^{a,b}$
Phe	2	$2^{a,c,h}$	Thr	7	$3^{a-c,e-g}$
Tyr	1	1 ^{e,h}	Lys	7	$1^{a,b,f,g}$
Glu	6	1 ^{b,c}	Ile	7	6^{a-g}

^a2D ¹⁵N-¹³C experiment on ¹³C2-ubiquitin.

In comparison, the amino acids produced from the citric acid cycle were more difficult to assign because of their similar labeling patterns, low labeling levels, and chemical shift overlap. Further experimental schemes, yielding larger ratios of chemical shift dispersion over linewidths, are necessary for tackling the assignment of these difficult amino acids.

Conclusions and outlook

The ability to assign the NMR resonances to the protein sequence provides the basis for efficient structure elucidation of extensively labeled proteins. We have shown here that multidimensional MAS correlation spectroscopy in conjunction with selective and extensive ¹³C labeling makes it possible to assign a significant fraction of the ¹³C and ¹⁵N resonances of a model protein in the solid state. The N(CO)CA technique vields sequential backbone connectivities, while 3D NCACB correlation substantially enhances the site resolution and permits amino acid type assignment. The proton-driven ¹³C spin diffusion experiment extends the type assignment by utilizing not only Ca and CB chemical shifts but also the chemical shifts of carbons further down the side chain. Finally, the ¹³C dipolar INADEQUATE experiment supplements the type-assignment by removing the dominant spectral diagonal, thus allowing the detection of low levels of spin pairs such as the $C\alpha$ and $C\beta$ of serine and the $C\beta$

and Cy of threonine. Combining all four experiments presented here, and the previous ¹⁵N-¹³C 2D correlation technique, we have achieved sequence-specific assignment (mainly in the backbone) for 25% of the amino acids in ubiquitin and type-assignment for 47% of the residues.

All correlation techniques demonstrated here utilize dipolar-mediated coherence transfer schemes, which allow more rapid transfer than scalar-couplingbased pulse sequences without undue loss of transfer specificity at short mixing times. In fact, one can take advantage of the mixing-time dependence of dipolar transfer to probe one- to multiple-bond connectivities. This has been shown to be useful in 2D ¹⁵N-¹³C correlation spectroscopy.

The assignment techniques shown here are ultimately limited by two factors: the spectral linewidths and the incomplete labeling of the selective and extensive ¹³C labeling protocol. Given the significant line narrowing already achieved by the selective and extensive ¹³C labeling approach, the resolution may be further enhanced by higher power ¹H decoupling, better shift dispersion under higher magnetic field strengths, more sophisticated sample preparation methods, and faster spinning. For the few directly bonded and highly labeled ¹³C spin pairs, such as Val Cα and Cβ, techniques that remove ¹³C-¹³C one-bond J-couplings will be useful for reducing the linewidths, although such experiments may be costly in sensitivity (Straus et al., 1996). Finally, more labeled samples with variations

b2D ¹⁵N-¹³C experiment on ¹³C1-ubiquitin. c2D N(CO)CA experiment on U-¹³C, ¹⁵N-ubiquitin.

d_{3D} NCACB experiment on ¹³C2-ubiquitin. e_{2D} ¹³C spin diffusion experiment on ¹³C2-ubiquitin.

^f2D ¹³C dipolar INADEQUATE experiment on ¹³C2-ubiquitin.

^g2D ¹³C dipolar INADEQUATE experiment on ¹³C1-ubiquitin.

^h1D ¹³C CP.

of the selective and extensive ¹³C labeling approach will also be valuable for obtaining additional assignment.

Even with the incomplete resonance assignment presented here, much structural information should now become accessible. For example, multiple backbone and side-chain torsion angles can be measured from a single experiment to provide information on the secondary structure as well as the local tertiary structure of proteins (Hong, 1999a). Further work along these lines is currently pursued in our laboratory.

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